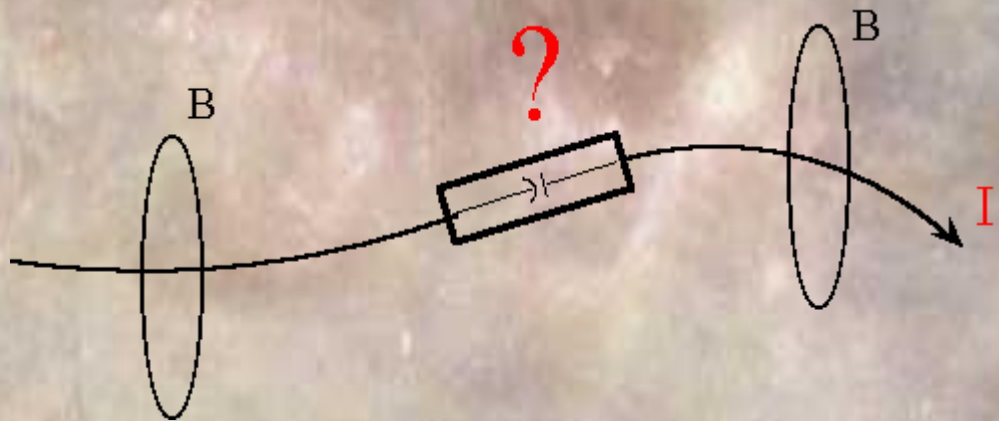
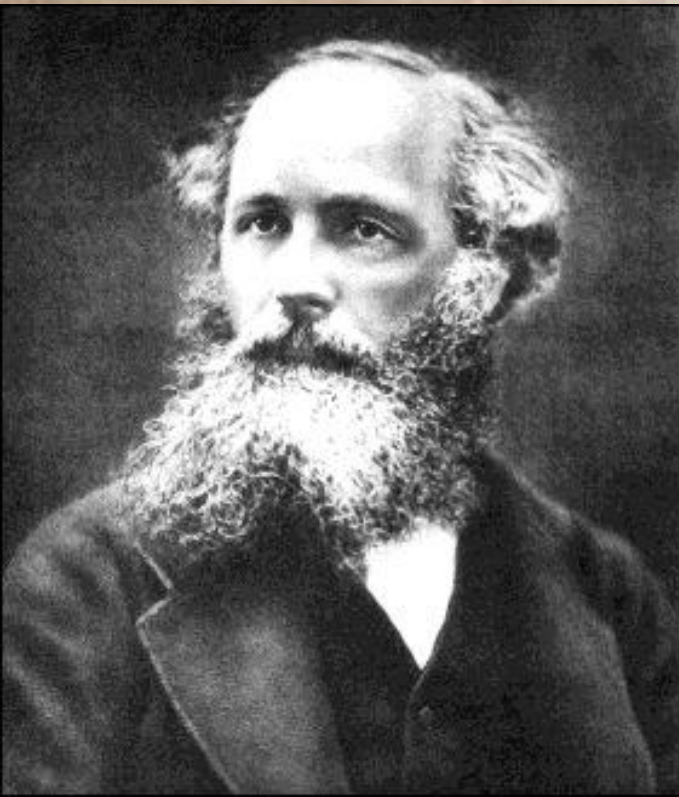


Maxwell's Equations and Electromagnetic Waves

Dr. Subhasis Chakrabarti

Setting the Stage - The Displacement Current

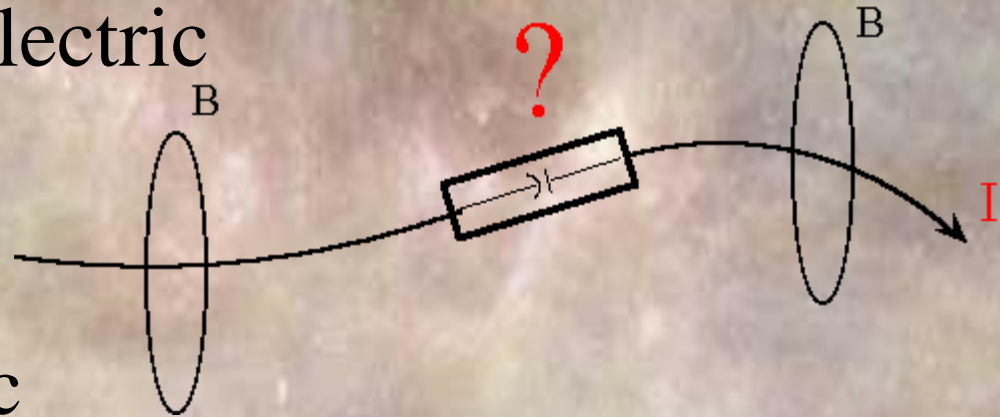
- Maxwell had a crucial “leap of insight”...



Will there still be a magnetic field around the capacitor?

A Beautiful Symmetry...

- A changing magnetic flux produces an Electric field



- A changing electric flux produces a Magnetic field

An extension of Ampere's Law...

- Maxwell reasoned that Ampere's Law would also apply to the **displacement current**.

$$\int B dl = \mu_o (I + I_{displacement})$$

Clever application of Gauss' Law here!

$$\epsilon_o \frac{d\phi_E}{dt} = \frac{dQ}{dt} = I_{displacement}$$

Maxwell's Equations (first glimpse)

- Faraday's Law:

$$\int \mathbf{E} d\mathbf{l} = -\frac{d}{dt} \int B_n dA$$

- Ampere's Law:

$$\int \mathbf{B} d\mathbf{l} = \mu_0 I + \mu_0 \epsilon_0 \frac{d}{dt} \int E_n dA$$

Maxwell's Equations – Integral Form

- Faraday's Law:
$$\int E dl = -\frac{d}{dt} \int B_n dA$$
- Ampere's Law:
$$\int B dl = \mu_o I + \mu_o \epsilon_o \frac{d}{dt} \int E_n dA$$
- Gauss' Law
$$\int E \cdot dA = \frac{Q}{\epsilon_o}$$
$$\int B \cdot dA = 0$$

Maxwell's Equations – Differential Form

$$\int E \cdot dA = \int_V (\nabla \cdot E) dV = \frac{Q}{\epsilon_0}$$

$$Q = \int_V \rho dV$$

$$\nabla \cdot E = \frac{\rho}{\epsilon_0}$$

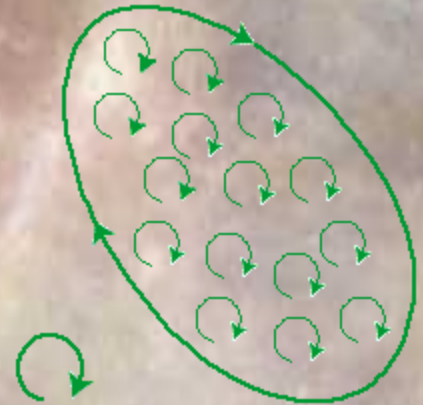
Gauss' Theorem - integral
of a flux equals volume
integral of divergence

Maxwell's Equations – Differential Form

$$\int E dl = -\frac{d}{dt} \int B_n dA$$

$$\int_{\Gamma} E \cdot dl = \int_{\Sigma} (\nabla \times E) dA = -\frac{d}{dt} \int B_n dA$$

$$\nabla \times E = -\frac{\partial B}{\partial t}$$



Stoke's Theorem: "Integral
around the path equals flux
of the curl"

Maxwell's Equations – Differential Form

$$\int B dl = \int_{\Sigma} (\nabla \times B) dA = \mu_0 \varepsilon_0 \frac{d}{dt} \int E_n dA$$

$$\nabla \times B = \mu_0 \varepsilon_0 \frac{\partial E}{\partial t}$$

Maxwell's Equations – Differential Form

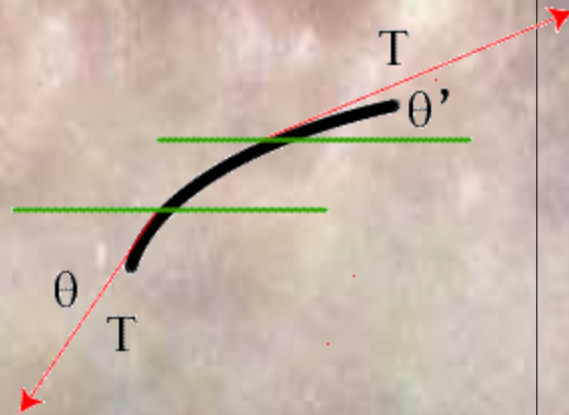
$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0} \quad \nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \cdot \mathbf{B} = 0 \quad \nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

The Wave Equation

Detour
Ahead

- How fast will a wave travel along a string of density ρ ?



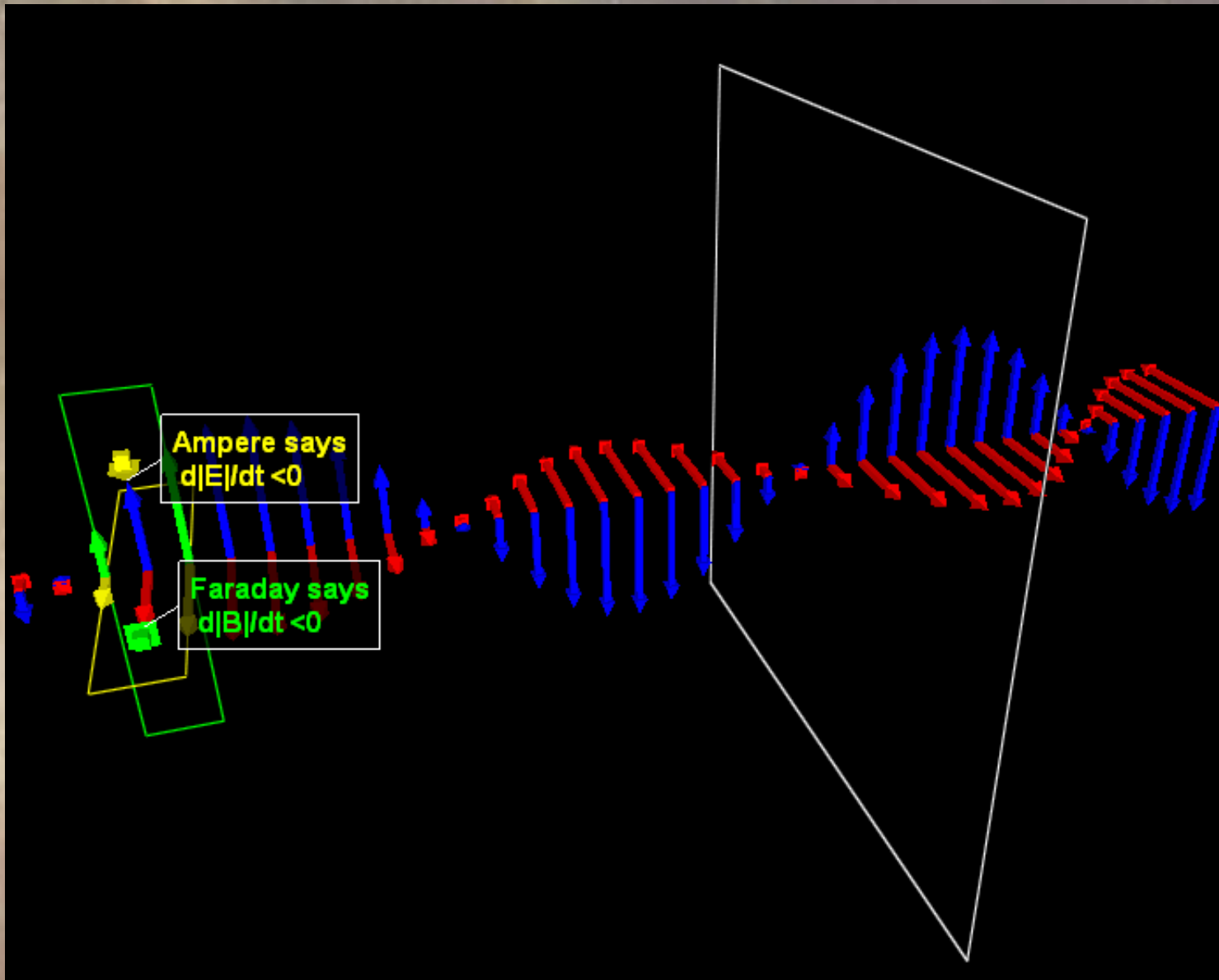
Two Ways to M'Es...

- Abstract: $\nabla \times (\nabla \times A) = \nabla(\nabla \cdot A) - \nabla^2 A$

- Physical:

- Imagine a plane wave of electric field in z-direction

$$E(x, t) = \begin{bmatrix} 0 \\ 0 \\ E_z(x, t) \end{bmatrix}$$



[Go to Rob Salgado's sim of this](#)

Moving Fields...

- Moving E-Field leads to...
- Moving B-Field leads to...

It's Alive!

- Well, at least it's a wave! Combining the last two equations leads us to:

$$\frac{\partial^2 E_y}{\partial x^2} = \mu_0 \epsilon_0 \frac{\partial^2 E_y}{\partial t^2}$$

- example - consider the electric field part of an electromagnetic wave described by:

$$E(x, t) = E_0 \sin(kx - \omega t) j + E_0 \cos(kx - \omega t) k$$

The Poynting Vector

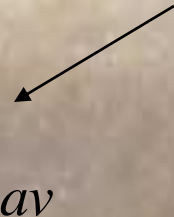
- Light waves (and all electromagnetic waves) carry energy

$$u = u_E + u_B = \frac{EB}{\mu_0 c}$$

- A wave has an intensity

$$I = u_{average} c$$

Poynting Vector

$$I = \frac{1}{2} \frac{E_o B_o}{\mu_0} = \left| \vec{S} \right|_{av}$$
$$\vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B}$$


Radiation Pressure

- Light waves (and all electromagnetic waves) exert pressure

$$P = \frac{I}{c}$$

Accelerating Charges Radiate Power!

- We can show (by dimensional argument) that an accelerating charge should radiate energy at a rate given by:

$$P = \alpha \frac{kq^2 a^2}{c^3}$$

- More detailed argumentation reveals that $\alpha = 2/3$